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Title: Method for determining the sound velocity in a base material, particularly for measuring the thickness of a wall

[0001] The invention relates to a method for determining the sound velocity C_b in a base material, using an ultrasonic probe comprising a transmit crystal, a receive crystal and a precursor body, said precursor body a) having a coupling surface by which the probe is couplable to the base material, b) receiving the receive crystal and the transmit crystal and c) having a sound velocity C_v , the transmit crystal and the receive crystal being oriented to be inclined towards each other and each towards the coupling surface so that a main transmission direction of the transmit crystal and a main receiving direction of the receive crystal intersect below the coupling surface, transmit crystal and receive crystal being spaced apart at a center to center distance K , the transmit crystal being spaced at a center to center distance D_s from the coupling surface and the receive crystal at a distance D_e from the coupling surface, by which method an ultrasonic pulse is generated by the transmit crystal, passes through the precursor body into the base material where it produces a creeping wave a portion of which reaches the receive crystal via the precursor body, and to a corresponding device.

[0002] A prerequisite of the ability of determining the wall thickness of the base material is to determine the sound velocity C_b . Although it is known to determine the wall thickness of a base material by multiple reflection of a pulse at an entrance surface and at a back surface of the base material, this method makes sufficiently mirror-like and, as a result thereof, smooth surfaces, more specifically a sufficiently smooth back surface a provision for ensuring multiple forward and reflected pulse waves within the base material. This method cannot be used with rough back surfaces; in this case, one is dependent on a pulse wave that travels back and forth only once. The wall thickness may then be determined through the sound velocity C_b .

[0003] A method and a device for determining the thickness of a coated base material is known from U.S. Patent No. 6,035,717. Using this method for determining the sound velocity C_b of the base material, the uncoated base material is first measured, a pulse is sent from a transmit crystal through the precursor body into the base material where a creeping wave is produced a portion of which being again decoupled and received by the receive crystal. Inasmuch, this patent coincides with the invention.

[0004] However, according to U.S. Patent No. 6,035,717, the path of the pulse is assumed to be imposed. It appears that the inventor of said U.S. Patent Document was well aware of the fact that assuming a geometrically imposed path along the main beams will result in certain inaccuracies in determining the sound velocity C_b . He therefore gives practical instructions, proposing to keep the distance separating the two crystals from the coupling surface as small as possible. This actually makes determining the sound velocity in the base material more precise, so that, put another way, the inaccuracy is reduced. However, a probe the precursor body of which has a short sound travel distance suffers from the disadvantage that but little material of the precursor body is available for the wear occurring during each practical testing so that the probe needs to be replaced earlier than a probe having a greater precursor distance.

[0005] This is where the invention comes in. Its object is to improve the method according to U.S. Patent No. 6,035,717 A so that the sound velocity C_b in the base material is determined with greater accuracy and so as to permit determining the thickness of a coating applied to said base material with greater accuracy as well and to allow using a probe the precursor distance of which is sufficiently thick.

[0006] In view of the features mentioned herein above, which are incorporated herein, this object is solved by measuring the shortest sound travel time T_{tot} and by determining the sound velocity C_b in the base material via that path that delivers, as a function of the sound velocity, the shortest total travel time T_{tot} between transmit crystal and receive crystal.

[0007] This method takes into account that the path by which the pulse travels through the precursor body, along the surface of the base material (as the surface wave) and back to the precursor body is influenced not only by the previously known variables K , D_v and C_v (K = center to center distance between the contact surfaces of the crystals, D_v = center to center distance between the contact surface of a crystal and the coupling surface) but also by the sound velocity C_b . If, compared to the sound velocity C_v in the precursor body, said sound velocity C_b is relatively high, the portion of the travel distance S_b along the surface of the base material will also be relatively large. If, by contrast, the sound velocity C_b in the base material is relatively small, the travel distance S_b of the surface wave in the base material will be relatively short, with the travel distances S_v within the precursor body becoming longer. Similar conditions will prevail when light is reflected between different optical media such as water and air. In this case again, the geometrically shortest path for a light pulse is not the shortest in time.

[0008] The achievement of the invention is that it relies on the observation that in determining the shortest travel time T_{tot} of the ultrasound pulse and in optimizing all of the possible sound travel paths leading to the sound travel path that supplies the shortest total travel time as a function of C_b , one obtains precise information about the sound velocity C_b in the base material. Accordingly, the invention relies on the actual paths travelled by a sound pulse. It does not, as this is the case with U.S. Patent No. 6,035,717, rely on any assumption about the path. Accordingly, the errors of this prior art

measurement method and of the related device are thus avoided in accordance with the invention.

[0009] Further advantages and characteristics of the invention will become apparent upon reviewing the other claims and the following non restrictive description of embodiments of the invention, given by way of example only with reference to the drawing, as well as of the method in accordance with the invention.

[0010] In said drawing:

[0011] FIG. 1: is a schematic representation as viewed from the side of a probe having two crystals, said probe being coupled to a base material, and shows the various portions of the overall path,

[0012] FIG. 2: is the representation according to FIG. 1 showing the travel distances, the sound velocities, and so on,

[0013] FIG. 3: is a representation like FIG. 1 with the base material having now additionally a thin layer (a coating) such as a color, a metal coating or a plastic layer and

[0014] FIG. 4: is a representation like FIG. 1 but now with two additional crystals for measuring the wall thickness.

[0015] The probe shown in FIG. 1 has a specially shaped, substantially prismatic precursor body 20. It has a level coupling surface 22, which is also referred to as the active surface, and, on the opposite side thereof, bevels on which a transmit crystal 24 and a receive crystal 26 are respectively retained, more specifically cemented. The two crystals 24, 26 are built according to the

same principle. They are disposed at an incline toward each other and also relative to the coupling surface 22. This arrangement will be discussed in further detail herein after.

[0016] A perpendicular bisector, meaning a line that perpendicularly bisects said crystal contact surface with the precursor body 20, is inclined at a certain angle relative to the coupling surface 22, said angle being $(90^\circ - \alpha_v)$ and being the same for the two crystals 24, 26. Further, the respective perpendicular bisectors lie in the same plane, namely in the plane of FIG. 1.

[0017] This may also be put another way: the two crystals 24, 26 exhibit two-fold symmetry about a plane of symmetry 32. They are inclined to the coupling surface 22 in such a manner that a surface wave 35 is generated in a base material 34 to which the precursor body 20 is coupled by well-known suited means, this being discussed in further detail herein after.

[0018] A separating layer 36 provided substantially along the plane of symmetry 32 prevents direct cross talk between transmit crystal 24 and receive crystal 26.

[0019] The perpendicular bisectors indicated usually coincide with a main beam, i.e., with a main transmitting beam 38 and with a main receiving beam 40.

[0020] The sound velocity C_v in the precursor body 20 is known. The distance K between the surface centers of the two crystals 24, 26 is known as well. Finally, the distance of the surface center of the transmit crystal 24 from the coupling surface 22 and the distance of the center of the receive crystal 26 from the coupling surface 22 can be determined and are thus known. Due to symmetry, they both have the value D_v . With only these data, it is now possible

to determine the sound velocity C_b in the base material 34. In a further step, the thickness, that is the wall thickness D_b of the base material 34, can be determined.

[0021] If the sound velocity C_b in the base material 34 is approximately as high as the sound velocity of steel, the shortest path of a sound pulse from the transmit crystal 24 to the receive crystal 26 is as follows: the pulse travels along the main transmitting beam, then in the form of a surface wave 35 in the base material 34 and eventually along the main receiving beam 40 into the receive crystal 26. This path is shown in a dashed line in FIG. 1; it extends along the main transmitting beam 38 and along the main receiving beam 40.

[0022] If however the sound velocity C_b in the base material 34 is smaller than that of steel, the sound path will use as much of a distance as possible within the precursor body 20, with the length of the distance S_b realized by the surface wave 38 in the base material 34, shortening. This case is shown in FIG. 1 in which the sound travel path 42 is shown in a dotted line.

[0023] If conversely the sound velocity C_b in the base material is higher than that of steel, the sound distance S_v within the precursor body 20 becomes shorter, with the travel distance S_b in the form of the surface wave 35 becoming longer as a result thereof. This case is represented in FIG. 1 which shows the sound travel path 44 in a dash-dot line.

[0024] For simplicity's sake, FIG. 1 only shows in a dashed line the complete sound travel path extending along the main beams 38, 40. It can be seen that the travel distance S_b of the surface wave 35 is a function of the sound velocity C_b in the base material 34 and moreover depends on the constant variables K , C_v and D_v . In accordance with the invention, the sound velocity C_b in the base material 34 is obtained by optimizing the associated sound travel path.

Accordingly, that sound travel path is taken as a basis that supplies the shortest total travel time T_{tot} as a function of the sound velocity C_b to be determined.

[0025] The greater the angle to the main beam, the more the amplitude of the sound pressure diminishes. If however one only measures the signal with the shortest total travel time T_{tot} , one is independent, within certain limits, of the amplitude of the receive signal. Ideally, the crystals 24, 26 should be spherical sources, but this is not the case. With the sound velocities obtained in practice, the fact that the crystals 24, 26 are not spherical has not such a great influence that this should be taken into consideration and specially evaluated. The orientation of the crystals 24, 26 ideally occurs for a mean value of the sound velocity C_b (e.g., for steel C_b about 6000 m/s).

[0026] Accordingly, the respective path by which the sound pulse with the shortest total travel time T_{tot} travels is function of the sound velocity C_b and further dependent on the known values K , D_v and C_v . If the structure does not exhibit two-fold symmetry, the various center-to-center distances of the transmit crystal 24 and of the receive crystal 26 from the coupling surface 22 must be taken into consideration.

[0027] According to FIG. 1, the ultrasound propagates from the transmit crystal 24 along a first path segment S_v until it reaches the base material 34, needing therefore the time T_v . There, a creeping wave 35 is generated. It has the length indicated in FIG. 2. It travels this length within time T_b . One portion of the creeping wave reaches the receive crystal 26 by a path which has, due to the symmetry, the length S_v and for which the time T_v is needed.

[0028] What is sought is the distance S_b between the starting point and the end point of the creeping wave or of the surface wave 35 respectively for longitudinal waves.

[0029] General formulas for sound propagation will be set forth herein after. Only the shortest total travel time T_{tot} will be considered. In order to determine the latter, it is necessary to take the total travel distance from the transmit crystal 24 to the receive crystal 26 into consideration. It will be assumed that K = center to center distance of the probe heads 24, 26; D_v = center to center distance of the crystals 24, 26 from the coupling surface 22 and C_v = sound velocity in the precursor body are known and time constant, for the short duration of the measurement at any rate.

[0030] An ultrasound pulse emitted by the transmit crystal 24 will not only effect a surface wave 35 in the base material 34, but other waves as well; the longitudinal surface wave 35 has the shortest travel distance and the shortest total travel time T_{tot} as well.

[0031] According to FIG. 2, the following applies:

[0032] (1) $S_b = K - 2K_b;$

[0033] (2) $\tan \alpha_v = \frac{K_b}{D_v}; \rightarrow K_b = D_v \tan \alpha_v;$

[0034] (3) $S_b = K - 2D_v \tan \alpha_v; \quad (1/2)$

[0035] Different sound travel paths differ by their entrance angle α_v .

[0036] This angle is maintained with assumption of the shortest possible total travel time T_{tot} :

[0037] (4) $T_{tot} = 2T_v + T_b;$

[0038] (5) $C_v = \frac{S_v}{T_v}; \rightarrow T_v = \frac{S_v}{C_v};$

[0039] (6) $C_b = \frac{S_b}{T_b}; \rightarrow T_b = \frac{S_b}{C_b};$

[0040] (7) $\cos \alpha_v = \frac{D_v}{S_v}; \rightarrow S_v = \frac{D_v}{\cos \alpha_v};$

[0041] (8) $T_v = \frac{D_v}{C_v \cos \alpha_v};$ (5/7)

[0042] (9) $T_{tot} = \frac{2D_v}{C_v \cos \alpha_v} + \frac{K - 2D_v \tan \alpha_v}{C_b};$ (4/8/3/6)

[0043] (10) $T_{tot} = \frac{K}{C_b} + 2D_v \left(\frac{1}{C_v \cos \alpha_v} - \frac{\tan \alpha_v}{C_b} \right);$ (9)

[0044] What is sought is the minimum of the linear function $T_{tot}(\alpha_v)$. This can for example be determined through the first derivation after the angle α_v ; the first derivation must be zero for a certain angle α_v and the second derivation must be positive:

[0045] (11) $\frac{\partial T_{tot}(\alpha_v)}{\partial \alpha_v} = 0!$

[0046] (12) $\frac{\partial T_{tot}(\alpha_v)}{\partial \alpha_v} = 2D_v \left(\frac{\sin \alpha_v}{C_v \cos^2 \alpha_v} - \frac{1}{C_b \cos^2 \alpha_v} \right) = 0;$ (10)

[0047] (13) $\frac{\sin \alpha_v}{C_v} - \frac{1}{C_b} = 0; \rightarrow \sin \alpha_v = \frac{C_v}{C_b}; \rightarrow \alpha_v = \arcsin\left(\frac{C_v}{C_b}\right);$ (12)

[0048] Taking now (3) into consideration, it will be recognized that the sound travel distance S_b is dependent on the two sound velocities C_v and C_b :

[0049] (14) $S_b = K - 2D_v \tan\left(\arcsin\left(\frac{C_v}{C_b}\right)\right);$ (3/12)

[0050] As it can be assumed that K , D_v and C_v are constant, this means that $S_b = f(C_b)$. The following equation (15) describes the dependency between the total travel time T_{tot} measured and the sound velocity C_b to be determined:

$$[0051] \quad (15) \quad T_{tot} = \frac{K}{C_b} + 2D_v \left(\frac{1}{C_v \cos(\arcsin(\frac{C_v}{C_b}))} - \frac{\tan(\arcsin(\frac{C_v}{C_b}))}{C_b} \right); \quad (9/12)$$

[0052] As a result, a clear correlation between the total travel time T_{tot} and the sound velocity C_b to be determined is achieved. All the other variables in the equation (15) are known and constant.

[0053] Accordingly, it is clearly possible to determine the sound velocity C_b in the base material from the total travel time T_{tot} . This finding now permits to determine the wall thickness D_b of the base material.

[0054] FIG. 3 shows the arrangement of the previous FIGS. with the base material 34 now having an additional layer 46, a so-called coating. This layer has a thickness D_s . It is to be determined from the sound velocity C_s of the layer. As the latter is not known either, the sound velocity is determined first like before.

[0055] FIG. 3 again only shows the sound travel distance with the shortest total travel time T_{tot} . Other propagations also occur, a surface wave is for example also generated at the surface of the coating 46 that is turned toward the coupling surface 22, but this wave is to occur in time after the surface wave 35 in the base material 34. This means that the sound velocity C_b in the base material 34 has to be sufficiently higher than the sound velocity C_s in the layer 46. In practice, this is mostly the case. Typically, the base material is a metal, the sound velocities range from 4500 to 7000 m/s. The layer 46 typically is a plastic material, a color or the like; the sound velocities typically range from 2000 to 3000 ms. In the event the sound velocity C_s in the layer 46 is relatively high, for example if the layer is a metal coating on a base material made from plastic material, the layer consists of a metal having a higher sound velocity

than the base material, e.g., if the coating is made from Ag and the base material from Au, it must be proceeded according to the above mentioned considerations, simply by utilizing the layer 46 as the base material.

[0056] Herein after, the sound travel path shown in FIG. 3 is considered the shortest travel path. In the layer 46, the entrance angle α_v is changed to α_s . The travel distances in the volume of the layer 46 are apparent from FIG. 3; they amount to S_s . The associated sound travel time is T_s .

[0057] For the shortest total travel time, the following then applies:

[0058] (16) $T_{tot} = 2(T_v + T_s) + T_b;$

[0059] (17) $C_v = \frac{S_v}{T_v}; \rightarrow T_v = \frac{S_v}{C_v};$

[0060] (18) $\cos \alpha_v = \frac{D_v}{S_v}; \rightarrow S_v = \frac{D_v}{\cos \alpha_v};$

[0061] (19) $T_v = \frac{D_v}{C_v \cos \alpha_v};$ (17/18)

[0062] In the layer 46, the conditions are the same; for the layer 46, the following applies:

[0063] (20) $T_s = \frac{D_s}{C_s \cos \alpha_s};$ (19)

[0064] (21) $T_b = \frac{S_b}{C_b};$

[0065] (22) $S_b = K - 2(K_b + K_s);$

[0066] (23) $\tan \alpha_v = \frac{K_s}{D_v}; \rightarrow K_s = D_v \tan \alpha_v;$

[0067] (24) $\tan \alpha_s = \frac{K_b}{D_s}; \rightarrow K_b = D_s \tan \alpha_s;$

[0068] (25) $S_b = K - 2(D_s \tan \alpha_s + D_v \tan \alpha_v);$ (22/23/24)

[0069] (26) $T_b = \frac{K - 2(D_s \tan \alpha_s + D_v \tan \alpha_v)}{C_b};$

(21/25)

[0070] Now all the elements for Ttot are known:

[0071] (27) $T_{tot} = 2\left(\frac{D_v}{C_v \cos \alpha_v} + \frac{D_s}{C_s \cos \alpha_s}\right) + \frac{K - 2(D_s \tan \alpha_s + D_v \tan \alpha_v)}{C_b};$

(16/19/20/26)

[0072] (28) $T_{tot} = \frac{2D_v}{C_v \cos \alpha_v} + \frac{2D_s}{C_s \cos \alpha_s} + \frac{K}{C_b} - \frac{2D_s \tan \alpha_s}{C_b} - \frac{2D_v \tan \alpha_v}{C_b};$ (27)

[0073] (29) $T_{tot} = \frac{K}{C_b} + 2\left(D_v\left(\frac{1}{C_v \cos \alpha_v} - \frac{\tan \alpha_v}{C_b}\right) + D_s\left(\frac{1}{C_s \cos \alpha_s} - \frac{\tan \alpha_s}{C_b}\right)\right);$ (28)

[0074] Now, the total travel time Ttot is not only a function of the entrance angle α_v (like in (15)), but also a function of the entrance angle α_s and can be represented as follows:

[0075] (30) $T_{tot}(\alpha_v, \alpha_s) = \frac{K}{C_b} + 2(f_1(\alpha_v) + f_2(\alpha_s));$ (29)

[0076] If the function Ttot(α_v , α_s) has a minimum this can again be established from the first derivation after the two angles. The first derivations must be 0:

[0077] (31) $\frac{\partial f_1(\alpha_v)}{\partial \alpha_v} = 0!$

[0078] (32) $\frac{\partial f_2(\alpha_s)}{\partial \alpha_s} = 0!$

[0079] (33) $\frac{\partial f_1(\alpha_v)}{\partial \alpha_v} = D_v \left(\frac{\sin \alpha_v}{C_v \cos^2 \alpha_v} - \frac{1}{C_b \cos^2 \alpha_v} \right) = 0$
(31/29)

[0080] (34) $\frac{\sin \alpha_v}{C_v} - \frac{1}{C_b} = 0; \rightarrow \frac{\sin \alpha_v}{C_v} = \frac{1}{C_b}; \rightarrow \sin \alpha_v = \frac{C_v}{C_b}; \rightarrow \alpha_v = \arcsin\left(\frac{C_v}{C_b}\right); (33)$

[0081] (35) $\frac{\partial f_2(\alpha_s)}{\partial \alpha_s} = D_s \left(\frac{\sin \alpha_s}{C_s \cos^2 \alpha_s} - \frac{1}{C_b \cos^2 \alpha_s} \right) = 0$
(32/29)

[0082] (36) $\frac{\sin \alpha_s}{C_s} - \frac{1}{C_b} = 0; \rightarrow \frac{\sin \alpha_s}{C_s} = \frac{1}{C_b}; \rightarrow \sin \alpha_s = \frac{C_s}{C_b}; \rightarrow \alpha_s = \arcsin\left(\frac{C_s}{C_b}\right); (35)$

[0083] The results (34) and (36) are now inserted into the equation (29) which yields:

[0084] (37)

$$T_{tot} = \frac{K}{C_b} + 2 \left(D_v \left(\frac{1}{C_v \cos \arcsin\left(\frac{C_v}{C_b}\right)} - \frac{\tan \arcsin\left(\frac{C_v}{C_b}\right)}{C_b} \right) + D_s \left(\frac{1}{C_s \cos \arcsin\left(\frac{C_s}{C_b}\right)} - \frac{\tan \arcsin\left(\frac{C_s}{C_b}\right)}{C_b} \right) \right);$$

(29/34/36))

[0085] The shortest total travel time Ttot can be measured using a suited measuring instrument, for example the apparatus DMS 2 of the applicant. From this total travel time, the thickness Ds of the layer 46 or the thickness Db of the base material 34 can now be determined, provided the sound velocities Cs and Cb are known. Cb for example can be determined at an uncoated site of the body being tested, using (27). Through an additional pair of crystals for thickness measurement with the crystals 48, which are built according to the same principle, the sound energy is emitted at a steep angle, as can be seen from FIG. 4. A back wall echo is generated the travel time of which is reduced by the travel time of Ds. The thickness Db of the base material is calculated therefrom with the sound velocity Cb being known.

[0086] A second echo, which occurs next in time, is obtained from a back wall 50 of the base material 34. The thickness D_b of the base material 34 can be determined from the time difference between the two echoes and from the previously measured sound velocity C_b in said base material. The thickness D_b can also be obtained as the difference between this echo of the back wall 50 and the entrance echo, taking into consideration the sound velocities and deducing the thickness D_s of the layer 46.

[0087] The following connections will become apparent or clear from the equation (37):

[0088] 1) $\frac{C_v}{C_b}$ and $\frac{C_s}{C_b}$ must be smaller than 1;9

[0089] 2) C_s can be smaller than C_v .